### **Assignment 1: selected solutions**

Problem 1: Show how the gates Y = -iZX,  $S = \sqrt{Z}$  and  $T = \sqrt{S}$  transform a general state of one quantum bit.

Solution: The gates *Z* and *X* transform the standard basis vectors as follows:

$$Z|0\rangle = |0\rangle, \quad Z|1\rangle = -|1\rangle,$$
  
 $X|0\rangle = |1\rangle, \quad X|1\rangle = |0\rangle,$ 

so the action of the gate Y = -iZX onto a general qubit state  $|\psi\rangle = c_0|0\rangle + c_1|1\rangle$  as

$$Y|\psi\rangle = -iZX(c_0|0\rangle + c_1|1\rangle)$$
  
=  $-iZ(c_0X|0\rangle + c_1X|1\rangle)$   
=  $-iZ(c_1|0\rangle + c_0|1\rangle)$   
=  $-i(c_1|0\rangle - c_0|1\rangle).$ 

From the following

$$\begin{split} Z|\psi\rangle &= Z(c_0|0\rangle + c_1|1\rangle) \\ &= c_0 Z|0\rangle + c_1 Z|1\rangle \\ &= c_0|0\rangle - c_1|1\rangle \\ &= c_0 S^2|0\rangle + c_1 S^2|1\rangle. \end{split}$$

we observe that  $S^2|0\rangle = |0\rangle = e^{i0}|0\rangle$  and  $S^2|1\rangle = -|1\rangle = e^{i\pi}|1\rangle$ . The action of the gate *S* onto a general state of one quantum bit

$$\begin{split} S |\psi\rangle &= S \left( c_0 |0\rangle + c_1 |1\rangle \right) \\ &= c_0 S |0\rangle + c_1 S |1\rangle \\ &= c_0 |0\rangle + e^{i\pi/2} c_1 |1\rangle \\ &= c_0 |0\rangle + i c_1 |1\rangle, \end{split}$$

and similarly

$$T|\psi\rangle = c_0|0\rangle + e^{i\pi/4}c_1|1\rangle.$$

# Single-qubit operations in the standard computational basis:

(i) Phase flip

$$\begin{aligned} \hat{Z} &= \left(\sum_{k=0,1} |k\rangle \langle k|\right) \hat{Z} \left(\sum_{l=0,1} |l\rangle \langle l|\right) = \sum_{k,l} \langle k|\hat{Z}|l\rangle |k\rangle \langle l| \\ &= \langle 0|\hat{Z}|0\rangle |0\rangle \langle 0| + \langle 0|\hat{Z}|1\rangle |0\rangle \langle 1| + \langle 1|\hat{Z}|0\rangle |1\rangle \langle 0| + \langle 1|\hat{Z}|1\rangle |1\rangle \langle 1| \\ &= \langle 0|\hat{Z}|0\rangle \left(\begin{array}{c} 1\\0\end{array}\right) \left(\begin{array}{c} 1&0\end{array}\right) + \langle 0|\hat{Z}|1\rangle \left(\begin{array}{c} 1\\0\end{array}\right) \left(\begin{array}{c} 0&1\end{array}\right) \\ &+ \langle 1|\hat{Z}|0\rangle \left(\begin{array}{c} 0\\1\end{array}\right) \left(\begin{array}{c} 1&0\end{array}\right) + \langle 1|\hat{Z}|1\rangle \left(\begin{array}{c} 0\\1\end{array}\right) \left(\begin{array}{c} 0&1\end{array}\right) \\ &= \left(\begin{array}{c} \langle 0|\hat{Z}|0\rangle & \langle 0|\hat{Z}|1\rangle \\ \langle 1|\hat{Z}|0\rangle & \langle 1|\hat{Z}|1\rangle\end{array}\right) = \left(\begin{array}{c} 1&0\\0&-1\end{array}\right) \end{aligned}$$

(ii) Bit flip

$$\hat{X} = \begin{pmatrix} \langle 0|\hat{X}|0\rangle & \langle 0|\hat{X}|1\rangle \\ \langle 1|\hat{X}|0\rangle & \langle 1|\hat{X}|1\rangle \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

Problem 2: Write the gates Y = -iZX,  $S = \sqrt{Z}$  and  $T = \sqrt{S}$  in the Dirac notation and in the matrix representation in the standard computational basis.

Solution: We can for example write the gate Z as  $Z = |0\rangle\langle 0| - |1\rangle\langle 1|$  which indeed transforms the standard basis states correctly, e.g.  $Z|0\rangle = |0\rangle\langle 0|0\rangle - |1\rangle\langle 1|0\rangle = |0\rangle$  and similarly  $X = |0\rangle\langle 1| + |1\rangle\langle 0|$ .

$$Y = -iZX$$
  
=  $-i(|0\rangle\langle 0| - |1\rangle\langle 1|)(|0\rangle\langle 1| + |1\rangle\langle 0|)$   
=  $-i|0\rangle\langle 1| + i|1\rangle\langle 0|$   
=  $-i\left(\begin{array}{cc} 1 & 0\\ 0 & -1\end{array}\right)\left(\begin{array}{cc} 0 & 1\\ 1 & 0\end{array}\right) = -i\left(\begin{array}{cc} 0 & 1\\ -1 & 0\end{array}\right) = \left(\begin{array}{cc} 0 & -i\\ i & 0\end{array}\right)$ 

For the rest of the problem 2 we use also functions of operators. After observing that the gate Z is diagonal in the standard basis, we get

$$S = \sqrt{Z} = \begin{pmatrix} \sqrt{1} & 0\\ 0 & \sqrt{-1} \end{pmatrix} = \begin{pmatrix} \sqrt{e^{i0}} & 0\\ 0 & \sqrt{e^{i\pi}} \end{pmatrix} = \begin{pmatrix} 1 & 0\\ 0 & i \end{pmatrix}$$

Problem 3: Show using the standard matrix representation that you can write the controlled-NOT operations using the projection operators as follows:

$$CNOT_{12} = \hat{P}_0 \otimes \hat{I} + \hat{P}_1 \otimes \hat{X}$$

where  $\hat{P}_0 = |0\rangle\langle 0|$  and  $\hat{P}_1 = |1\rangle\langle 1|$  are the single-qubit projectors, and  $\hat{I}$  is the single-qubit identity, i.e.  $2 \times 2$  unit matrix, and  $\hat{X}$  is the single-qubit bit-flip operator.

Solution: We first have to construct the matrix representation. The gate  $CNOT_{12}$  acts on the standard basis elements as  $CNOT_{12}|00\rangle = |00\rangle$ ,  $CNOT_{12}|01\rangle = |01\rangle$ ,  $CNOT_{12}|10\rangle = |11\rangle$  and  $CNOT_{12}|11\rangle = |10\rangle$ , which gives its matrix representation

$$CNOT_{12} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

We have to show that the r.h.s. gives the same matrix

Problem 4: Consider the controlled-*PHASE* gate, that flips the phase of the target qubit if the controlled qubit is in the state  $|1\rangle$ , and show in the standard matrix representation that  $CPHASE_{12} = CPHASE_{21}$ .

Solution:

and

$$\begin{aligned} CPHASE_{21} &= \hat{I} \otimes \hat{P}_0 + \hat{Z} \otimes \hat{P}_1 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \otimes \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} + \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \otimes \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} \\ &= \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} + \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \end{aligned}$$

which shows that

 $CPHASE_{12} = CPHASE_{21}.$ 

Problem 5: Consider the controlled-S gates, that applies the *S* gate to the target qubit if the controlled qubit is in the state  $|1\rangle$ . Write down the controlled-S gates  $CS_{12}$  and  $CS_{21}$  (i) in the standard matrix representation and (ii) in compact form using appropriate projectors.

Solution:

and similarly

$$CS_{21} = \hat{I} \otimes \hat{P}_0 + \hat{S} \otimes \hat{P}_1 = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & i \end{pmatrix}.$$

### **Assignment 2: supporting material**

## **Types of operators**

- 1. The hermitian conjugate or adjoint of a matrix is given as  $\hat{A}^{\dagger} = (\hat{A}^T)^*$ .
- 2. An operator  $\hat{A}$  is called **hermitian** if  $\hat{A}^{\dagger} = \hat{A}$ , or  $\langle \hat{A}\phi | \psi \rangle = \langle \phi | \hat{A}\psi \rangle$ .
- 3. An operator  $\hat{U}$  is called **unitary** if  $\hat{U}^{\dagger} = \hat{U}^{-1}$ , that is  $\hat{U}\hat{U}^{\dagger} = \hat{U}^{\dagger}\hat{U} = \hat{1}$ .
- 4. An operator  $\hat{P}$  satisfying  $\hat{P} = \hat{P}^{\dagger}$  and  $\hat{P} = \hat{P}^2$  is a **projection operator** or **projector**.

Example:

$$1.\ CNOT_{12}^{\dagger} = (CNOT_{12}^{T})^{*} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$$

2.  $CNOT_{12} = CNOT_{12}^{\dagger}$  so the operator is hermitian. The matrix representation of  $CNOT_{12}$  is real and symmetric.

3.  $(CNOT_{12})(CNOT_{12})^{\dagger} = (CNOT_{12})^{\dagger}(CNOT_{12}) = (CNOT_{12})(CNOT_{12}) = \hat{I}$ , so the operation is unitary.

4.  $CNOT_{12}$  is hermitian but it is not idempotent  $(CNOT_{12})(CNOT_{12}) \neq (CNOT_{12})$ , so it is not a projector.

### **Density operator/matrix**

We can represent a qubit state  $|\phi\rangle$ , and any quantum state, by the projector onto the one-dimensional subspace it spans:

$$\hat{\rho} = |\phi\rangle\langle\phi| = (c_0|0\rangle + c_1|1\rangle) \left(c_0^*\langle0| + c_1^*\langle1|\right) \\ = |c_0|^2 |0\rangle\langle0| + c_0c_1^*|0\rangle\langle1| + c_0^*c_1|1\rangle\langle0| + |c_1|^2 |1\rangle\langle1|$$

In matrix representation given by the standard computational basis, we have

$$\hat{\rho} = \begin{pmatrix} \rho_{00} & \rho_{01} \\ \rho_{10} & \rho_{11} \end{pmatrix} = \begin{pmatrix} c_0 \\ c_1 \end{pmatrix} \begin{pmatrix} c_0^* & c_1^* \end{pmatrix} = \begin{pmatrix} |c_0|^2 & c_0c_1^* \\ & & \\ c_0^*c_1 & |c_1|^2 \end{pmatrix}$$

We observe that the **norm** of a state is  $Tr(\hat{\rho}) = |c_0|^2 + |c_1|^2 = 1$  and also that  $\rho_{10} = \rho_{01}^*$ .

### **Bloch representation**

The single-qubit density matrix can be decomposed as follows

$$\hat{\rho} = \frac{1}{2} (\hat{I} + \vec{r} \cdot \vec{\sigma}) = \frac{1}{2} (\hat{I} + r_x \, \sigma_x + r_y \, \sigma_y + r_z \, \sigma_z)$$

$$= \frac{1}{2} \left[ \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + r_x \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} + r_y \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} + r_z \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \right]$$

$$= \frac{1}{2} \begin{pmatrix} 1 + r_z & r_x - ir_y \\ r_x + ir_y & 1 - r_z \end{pmatrix}$$

where  $\sigma_x$ ,  $\sigma_y$ , and  $\sigma_z$  are the Pauli matrices and the vector  $\vec{r} = (r_x, r_y, r_z)$  is called the **Bloch vector**.

The Bloch vector components are real numbers between 0 between 1, are related to the density matrix elements as shown (for a pure state)

$$r_x = 2 \operatorname{Re}(\rho_{10}) = 2 \operatorname{Re}(c_0^* c_1)$$
  

$$r_y = 2 \operatorname{Im}(\rho_{10}) = 2 \operatorname{Im}(c_0^* c_1) .$$
  

$$r_z = \rho_{00} - \rho_{11} = |c_0|^2 - |c_1|^2$$

1. 
$$|\phi\rangle = |0\rangle$$
  
 $\hat{\rho} = |0\rangle\langle 0| = \begin{pmatrix} 1\\0 \end{pmatrix} \begin{pmatrix} 1^* & 0^* \end{pmatrix} = \begin{pmatrix} 1 & 0\\0 & 0 \end{pmatrix} \implies \vec{r} = (0, 0, 1)$ 

2. 
$$|\phi\rangle = |1\rangle$$
  
 $\hat{\rho} = |1\rangle\langle 1| = \begin{pmatrix} 0\\1 \end{pmatrix} \begin{pmatrix} 0^* & 1^* \end{pmatrix} = \begin{pmatrix} 0 & 0\\0 & 1 \end{pmatrix} \implies \vec{r} = (0, 0, -1)$ 

3. 
$$|\phi\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$$
  

$$\hat{\rho} = |\phi\rangle\langle\phi| = \begin{pmatrix}\frac{1}{\sqrt{2}}\\\frac{1}{\sqrt{2}}\end{pmatrix}\begin{pmatrix}\frac{1}{\sqrt{2}}^* & \frac{1}{\sqrt{2}}^*\end{pmatrix} = \frac{1}{2}\begin{pmatrix}1 & 1\\\\1 & 1\end{pmatrix} \Rightarrow \vec{r} = (1, 0, 0)$$

4. 
$$|\phi\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$$
  
 $\hat{\rho} = |\phi\rangle\langle\phi| = \begin{pmatrix} \frac{1}{\sqrt{2}} \\ -\frac{1}{\sqrt{2}} \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{2}}^* & -\frac{1}{\sqrt{2}}^* \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1 & -1 \\ -1 & 1 \end{pmatrix} \Rightarrow \vec{r} = (-1, 0, 0)$ 

5. 
$$|\phi\rangle = \frac{1}{\sqrt{2}}(|0\rangle + i|1\rangle)$$
  
 $\hat{\rho} = |\phi\rangle\langle\phi| = \begin{pmatrix} \frac{1}{\sqrt{2}} \\ i\frac{1}{\sqrt{2}} \end{pmatrix} \left( \frac{1}{\sqrt{2}}^* \left(i\frac{1}{\sqrt{2}}\right)^* \right) = \frac{1}{2} \begin{pmatrix} 1 & -i \\ & \\ i & 1 \end{pmatrix} \Rightarrow \vec{r} = (0, 1, 0)$